

How ψ -ontic are ψ -ontic models? A case study

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Rose is a Rose is a Rose

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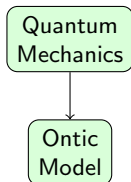
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A student bursts into the study of his professor and calls out: "Dear professor, dear professor! I have discovered a perpetual motion of the second kind!" The professor scarcely takes his eyes of his book and curtly replies: "Come back when you have found a neighborhood U of a state x_0 of such a kind that every $x \in U$ is connected with x_0 by an adiabat." – Walter

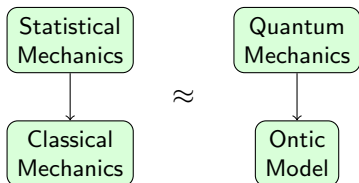
- 1 Ontic models framework.
- 2 When is an ontic model ψ -ontic?
- 3 The Kochen-Specker Theorem.
- 4 The ontic models of Meyer, Clifton and Kent (MKC).
- 5 The MKC models are ψ -ontic.
- 6 The MKC models are not very ψ -ontic.

Ontic models as complete descriptions



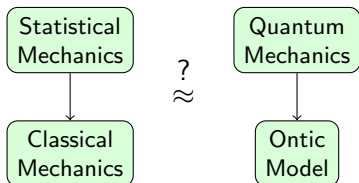
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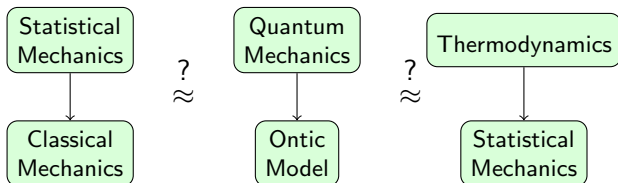
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- Do all phenomena described by SM also have a purely classical mechanical description?
- Theorems that constraint ontic models characterize the possible reducing theories for quantum mechanics.

Operational model $(\mathcal{P}, \mathcal{M})$

$P \in \mathcal{P}$ is a preparation, $M \in \mathcal{M}$ is a measurement,

$\mathbb{P}(m|M, P)$ probability of outcome m for measurement M performed on a system prepared according to P .

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Quantum mechanics as an operational model:

\mathcal{P} Set of quantum states,

\mathcal{M} Set of self-adjoint operators,

$$\mathbb{P}(a|A, |\psi\rangle) = \langle \psi | P_a^A | \psi \rangle.$$

Ontic model (Λ, Π, Ξ)

$\lambda \in \Lambda$ is an ontic state, $\xi \in \Xi$ is a Markov kernel
 $\mu \in \Pi$ is a probability measure, from Λ to Ω_ξ .

$\forall (P, M)$ there exists (μ_P, ξ_M) such that

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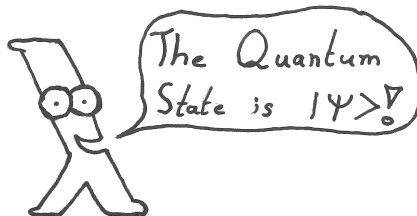
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Theorem: An ontic model always exists

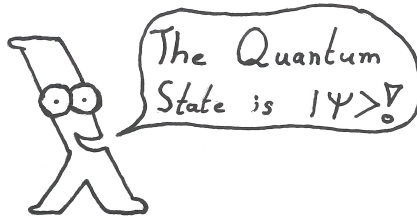
“Proof”:

$$\Lambda = \mathcal{P}, \mu_P(\{P'\}) = \delta_{PP'}, \xi_M(m|P) = \mathbb{P}(m|M, P).$$

The ontic state tells us what the quantum state is

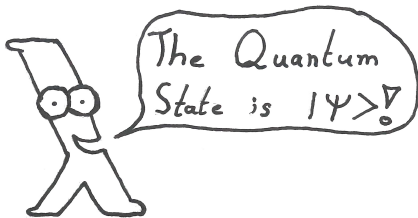


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- There exists a function $f : \Lambda \rightarrow \mathcal{P}$
 $f(\lambda)$ is the true quantum state of the system,
(bridge law as identity relation),
- Compatibility with quantum mechanical notion of quantum states:

$$\mu_{\psi}(f^{-1}(|\psi\rangle)) = 1.$$

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- 2 If $|\psi\rangle \neq |\phi\rangle$, then μ_ψ and μ_ϕ are non-overlapping.

Definition

An ontic model is ψ -ontic if for all $|\psi\rangle, |\phi\rangle \in \mathcal{P}$

$$D(\mu_\psi, \mu_\phi) = \sup_{\Omega \in \Sigma} |\mu_\psi(\Omega) - \mu_\phi(\Omega)| = 1.$$

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In statistical mechanics:

- If \mathcal{P} is the set of micro-canonical ensembles, then classical mechanics is P -ontic.
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Is just the energy ontic, or also the associated micro-canonical ensemble?

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Case study: The MKC hidden variable models.

Ontic Models and Hidden Variables

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The MKC models are constructed precisely to show the logical possibility of a **non-contextual hidden variable theory**. Allegedly, this possibility was ruled out by the Kochen-Specker theorem.

Assigning values to observables

For every non-degenerate self-adjoint operator A , there is a unique orthonormal basis

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such that

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For every non-degenerate self-adjoint operator A and every function f , the observables $f(A)$ and A can be measured jointly and the outcome for $f(A)$ is $f(a_j)$.

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For every non-degenerate self-adjoint operator A and every function f , the observables $f(A)$ and A can be measured jointly and the outcome for $f(A)$ is $f(a_j)$.

- The definite value assigned to $f(A)$ is determined by the definite value assigned to A .

Non-contextuality of quantum states

For two non-degenerate self-adjoint operators A, B , if

$$A|\psi\rangle = a|\psi\rangle \quad \text{and} \quad B|\psi\rangle = b|\psi\rangle$$

then for every quantum state $|\phi\rangle$

$$\mathbb{P}_{|\phi\rangle}(A = a) = \mathbb{P}_{|\phi\rangle}(B = b) = |\langle\phi|\psi\rangle|^2.$$

Defending non-contextuality

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Non-contextuality of definite values

For two non-degenerate self-adjoint operators A, B , if

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then A has the value a iff B has the value b .

The Kochen-Specker Theorem

Assigning definite values

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For every ontic state: $|e\rangle$ is the “true” vector in an orthonormal basis \mathcal{B} iff it is the true vector in every other orthonormal basis \mathcal{B}' that contains $|e\rangle$.

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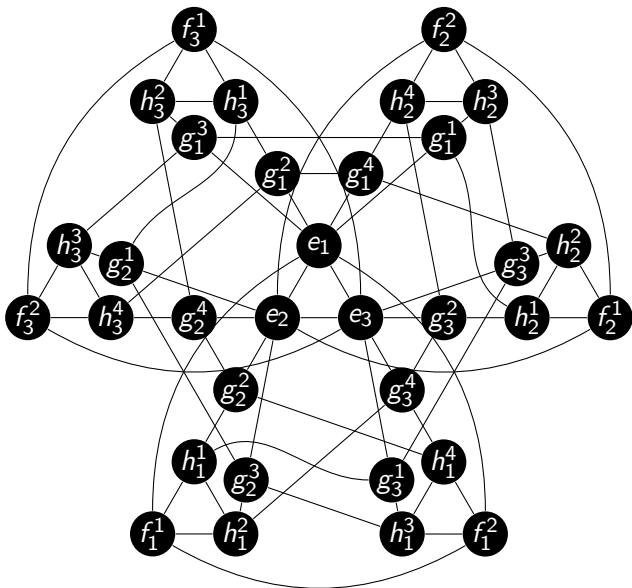
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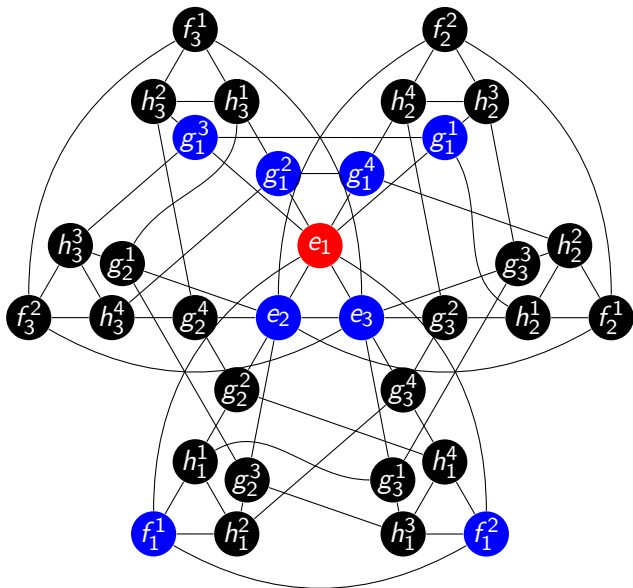
Kochen-Specker Theorem

There is a finite set of orthonormal bases for which one cannot select true vectors in a non-contextual way.

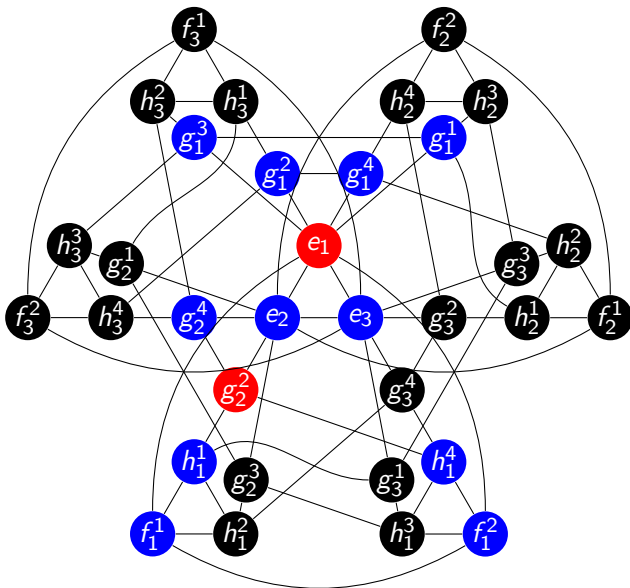
Uncolorable graph in 3 dimensions (Peres-cube)



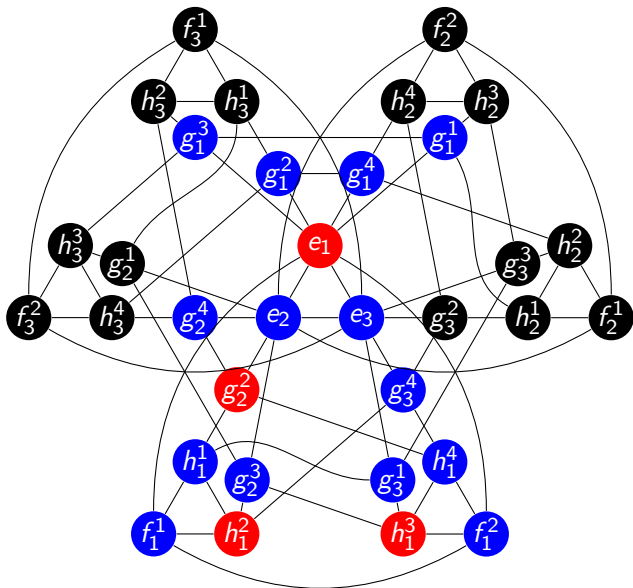
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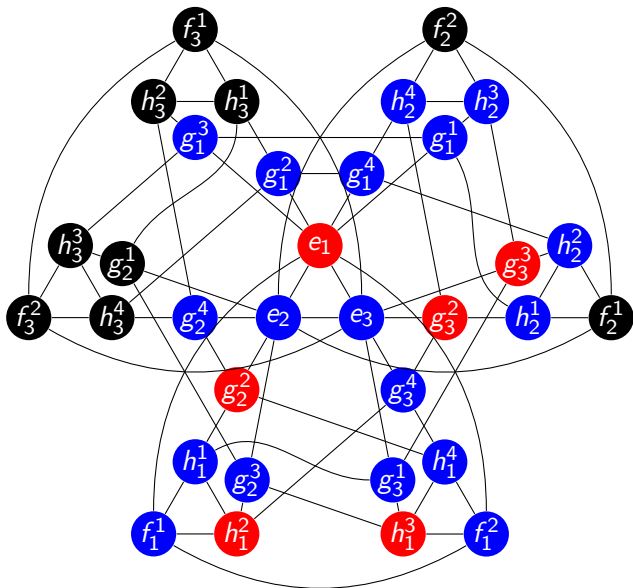
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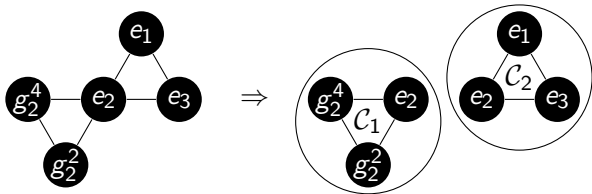
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Circumventing the Kochen-Specker Theorem

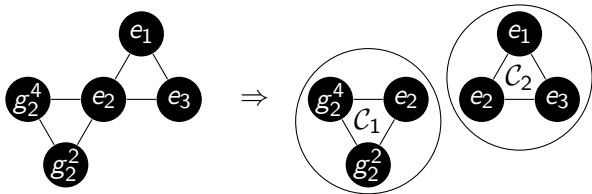
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Bell: Value assigned to an observable depends on the context \mathcal{C}

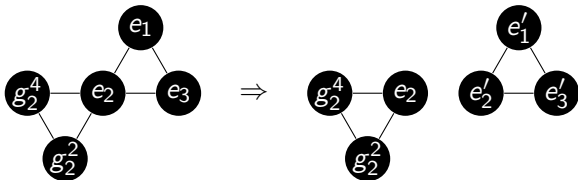


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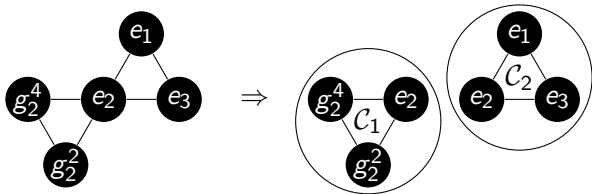


MKC: Not every orthonormal basis represents an observable

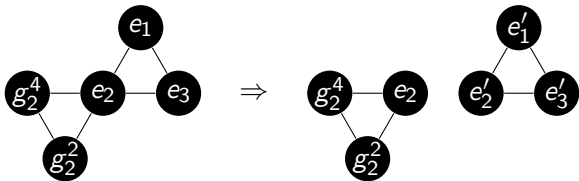


Circumventing the Kochen-Specker Theorem

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MKC: Not every orthonormal basis represents an observable



but every orthonormal basis can be approximated by an observable

$$0 < \|e_2 - e'_2\| < \epsilon.$$

The ontic models of Meyer, Kent and Clifton (MKC)

Definition

Two orthonormal bases

$$\mathcal{B}_1 = \{|e_1^1\rangle, \dots, |e_n^1\rangle\}, \quad \mathcal{B}_2 = \{|e_1^2\rangle, \dots, |e_n^2\rangle\}$$

are *totally incompatible* if $\forall i, j = 1, \dots, n$:

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Trivial Theorem

Let \mathfrak{B} be any set of pairwise totally incompatible orthonormal bases. The set of all self-adjoint operators with eigenvectors in one of the bases in \mathfrak{B} is colorable.

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Non-Trivial Theorem (Clifton & Kent)

There exists a countable set \mathfrak{B} of pairwise totally incompatible orthonormal bases that lies dense in the set of all orthonormal bases.

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- $|e_{i_1}^{k_1}\rangle \in \mathcal{B}_{k_1}, |e_{i_2}^{k_2}\rangle \in \mathcal{B}_{k_2}, \dots, |e_{i_n}^{k_n}\rangle \in \mathcal{B}_{k_n}$:

$$C_{e_{i_1}^{k_1}, \dots, e_{i_n}^{k_n}} := \bigcap_{j=1}^n \{\lambda \in \Lambda \mid \lambda(k_j) = i_j\},$$
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BCLM Theorem

Any ontic model that reproduces the predictions of QM is not maximally ψ -epistemic (but “almost” ψ -ontic).

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For cylinder sets:

$$\mu_\psi(C_{e_{i_1}^{k_1}, \dots, e_{i_n}^{k_n}}) := \prod_{j=1}^n |\langle \psi | e_{i_j}^{k_j} \rangle|^2 \rightarrow 1,$$

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Are the MKC models ψ -ontic?

ψ -onticness:

For all pairs of quantum states ψ, ϕ , for all corresponding probability measures $\mu \in \Delta_\psi, \nu \in \Delta_\phi$ in the ontic model, the variational distance $D(\mu, \nu) := \sup_{\Omega \in \Sigma} |\mu(\Omega) - \nu(\Omega)|$ equals 1.

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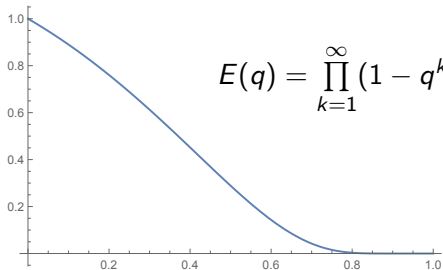
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Finding ontic states for $|\psi\rangle$

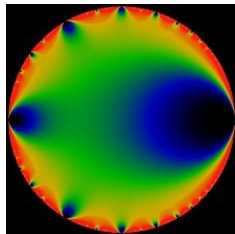
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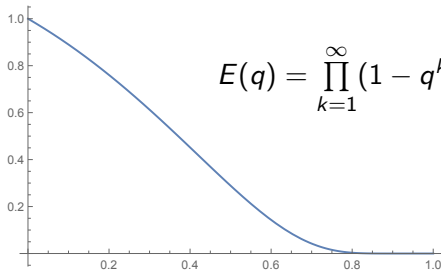


$$E(q) = \prod_{k=1}^{\infty} (1 - q^k)$$

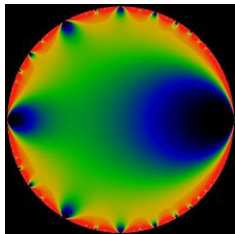


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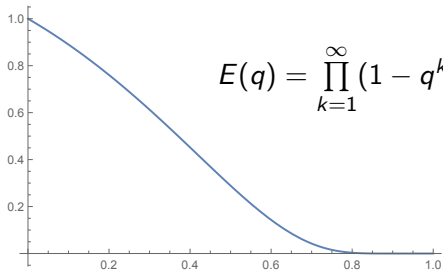
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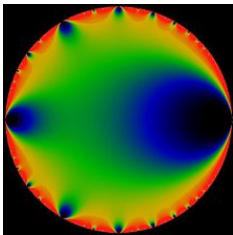
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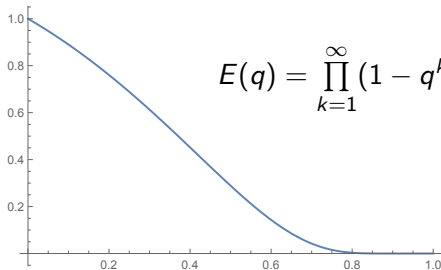
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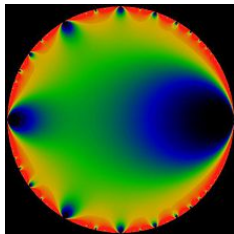
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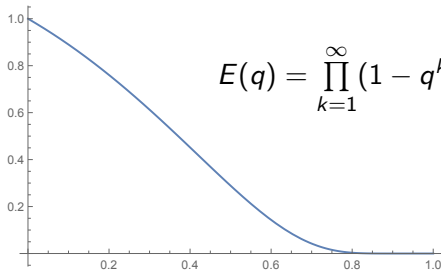
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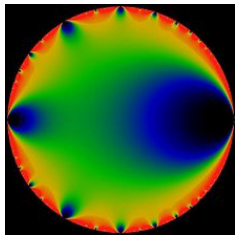
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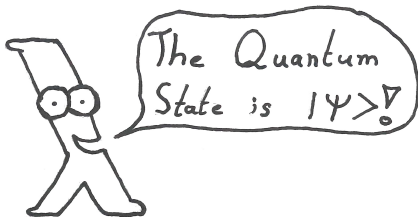
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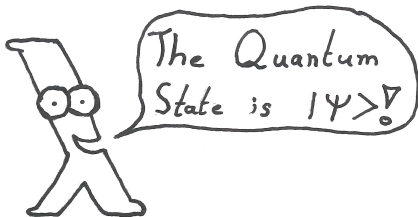
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How ψ -ontic are the ψ -ontic MKC models?



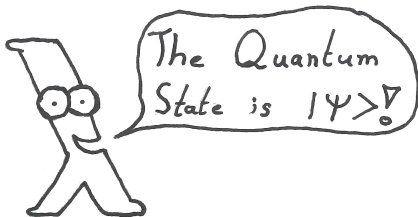
- 1 There exists a function $f : \Lambda \rightarrow \mathcal{P}$, compatible with QM: $\mu_{|\psi\rangle}(f^{-1}(|\psi\rangle)) = 1$.
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 - But we want the ontic state to tell us what $|\psi\rangle$ is, i.e., $\Lambda_\psi \subset \Lambda$ such that

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- The set Λ_ψ^ϵ is underdetermined by (ψ, ϵ) .
- More generally: $f : \Lambda \rightarrow \mathcal{P}$ need not be unique, so we cannot speak of an identity relation.