$\psi$ -onticity in the PBR theorem A case study using the MKC nullification of the KS theorem

Philosophy of Physics Seminar 8 May 2018

Ronnie Hermens

### Outline

- The PBR Theorem
- ullet Two notions of  $\psi$ -onticity
- The KS Theorem
- The MKC models
- ullet  $\psi$ -onticity of the MKC models

### The PBR Theorem



#### Formal statement

Any ontic model that reproduces the predictions of QM and satisfies the Preparation Independence Postulate is  $\psi$ -ontic.

### As a slogan

The quantum state is real.

# Methodology

#### PBR Theorem

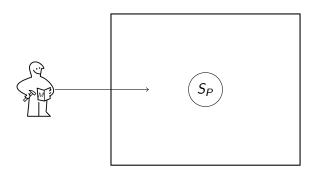
Any ontic model that reproduces the predictions of QM and satisfies the Preparation Independence Postulate is  $\psi$ -ontic.

- Quantum mechanics is viewed as an operational theory.
  - Quantum states are viewed as preparation procedures.
  - Quantum observables are viewed as measurement procedures.

## The operational approach

### Operational Prepare Measurement model $(\mathcal{P},\mathcal{M})$

 $P \in \mathcal{P}$  is a preparation,  $M \in \mathcal{M}$  is a measurement,  $\mathbb{P}(m|M,P)$  probability of outcome m for measurement M performed on a system prepared according to P.





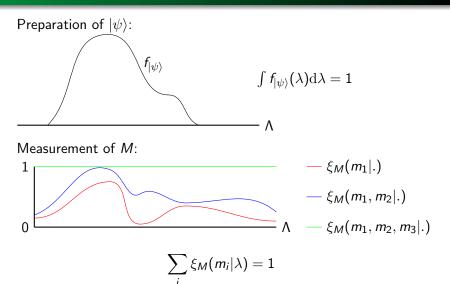
## Methodology

#### PBR Theorem

Any ontic model that reproduces the predictions of QM and satisfies the Preparation Independence Postulate is  $\psi$ -ontic.

- Quantum mechanics is viewed as an operational theory.
  - Quantum states are viewed as preparation procedures.
  - Quantum observables are viewed as measurement procedures.
- The ontic models framework is a general framework in which systems are assigned states.
  - State space Λ.
  - Preparations are identified with probability distributions over  $\Lambda$ .
  - Measurements are identified with response functions.

### Ontic models



Compatibility: 
$$\mathbb{P}(m|M,P) = \int \xi_M(m|\lambda) f_{|\psi\rangle}(\lambda) d\lambda$$
.

# Methodology

#### PBR Theorem

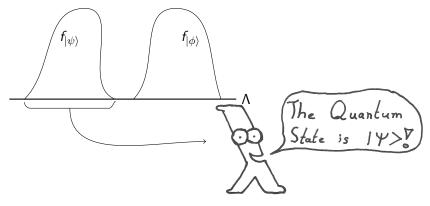
Any ontic model that reproduces the predictions of QM and satisfies the Preparation Independence Postulate is  $\psi$ -ontic.

- Quantum mechanics is viewed as an operational theory.
  - Quantum states are viewed as preparation procedures.
  - Quantum observables are viewed as measurement procedures.
- The ontic models framework is a general framework in which systems are assigned states.
  - State space Λ.
  - ullet Preparations are identified with probability distributions over  $\Lambda$ .
  - Measurements are identified with response functions.
- The ontic state  $\lambda$  determines the quantum state  $|\psi\rangle$ .

## $\psi$ -onticity

The ontic state  $\lambda$  determines the quantum state  $|\psi\rangle$ :

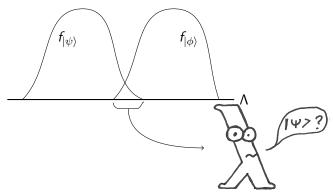
An important step towards the derivation of our result is the idea that the quantum state is physical if distinct quantum states correspond to non-overlapping distributions for  $\lambda$ .



## $\psi$ -onticity

The ontic state  $\lambda$  determines the quantum state  $|\psi\rangle$ :

An important step towards the derivation of our result is the idea that the quantum state is physical if distinct quantum states correspond to non-overlapping distributions for  $\lambda$ .



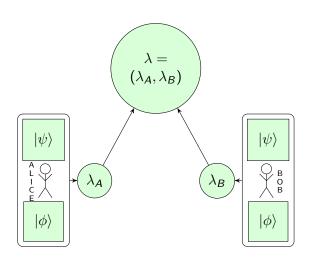
# Methodology

#### PBR Theorem

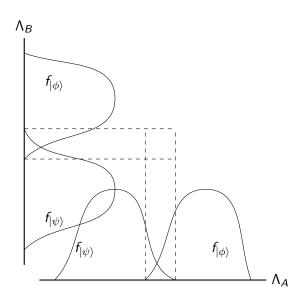
Any ontic model that reproduces the predictions of QM and satisfies the Preparation Independence Postulate is  $\psi$ -ontic.

- Quantum mechanics is viewed as an operational theory.
  - Quantum states are viewed as preparation procedures.
  - Quantum observables are viewed as measurement procedures.
- The ontic models framework is a general framework in which systems are assigned states.
  - State space Λ.
  - Preparations are identified with probability distributions over  $\Lambda$ .
  - Measurements are identified with response functions.
- The ontic state  $\lambda$  determines the quantum state  $|\psi\rangle$ .

# Preparation Independence



# Preparation Independence



### Without PIP

#### PBR Theorem

Any ontic model that reproduces the predictions of QM and satisfies the Preparation Independence Postulate is  $\psi$ -ontic.

#### **BCLM Theorem**

Any ontic model that reproduces the predictions of QM is almost  $\psi$ -ontic.

### Outline

- The PBR Theorem
- ullet Two notions of  $\psi$ -onticity
- The KS Theorem
- The MKC models
- ullet  $\psi$ -onticity of the MKC models

## $\psi$ -ontic Ontic Models

The ontic state tells us what the quantum state is



- There exists a function  $f:\Lambda o \mathcal{P}$   $f(\lambda)$  is the true quantum state of the system,
- Compatibility with quantum mechanical notion of quantum states:

$$\mu_{\psi}(f^{-1}(|\psi\rangle)) = 1.$$

## $\psi$ -ontic Ontic Models

Pusey, Barrett, Rudolph:

An important step towards the derivation of our result is the idea that the quantum state is physical if distinct quantum states correspond to non-overlapping distributions for

② If  $|\psi\rangle \neq |\phi\rangle$ , then  $\mu_{\psi}$  and  $\mu_{\phi}$  are non-overlapping.

#### Definition

An ontic model is  $\psi$ -ontic if for all  $|\psi\rangle$ ,  $|\phi\rangle \in \mathcal{P}$ 

$$D(\mu_{\psi},\mu_{\phi}) = \sup_{\Omega \in \Sigma} |\mu_{\psi}(\Omega) - \mu_{\phi}(\Omega)| = 1.$$

## Are $\psi$ -ontic Ontic Models $\psi$ -ontic?

- 2 If  $|\psi\rangle \neq |\phi\rangle$ , then  $\mu_{\psi}$  and  $\mu_{\phi}$  are non-overlapping.

#### Definition

An ontic model is  $\psi$ -ontic if for all  $|\psi\rangle\,, |\phi\rangle\in\mathcal{P}$ 

$$D(\mu_{\psi}, \mu_{\phi}) = \sup_{\Omega \in \Sigma} |\mu_{\psi}(\Omega) - \mu_{\phi}(\Omega)| = 1.$$

- Clearly  $1 \Rightarrow 2$ , but does  $2 \Rightarrow 1$ ?
- Intuitively: take  $\Omega = \Lambda_{\psi} := f^{-1}(|\psi\rangle)$ .
- But f need not exist for  $\psi$ -ontic models.

Case study: The MKC hidden variable models.

### Outline

- The PBR Theorem
- ullet Two notions of  $\psi$ -onticity
- The KS Theorem
- The MKC models
- ullet  $\psi$ -onticity of the MKC models

### Ontic Models and Hidden Variables

 Hidden variable models are traditionally concerned with the question:

Do measurements simply reveal the value of an observable, or is this value in some sense 'created' by the act of measurement?

- Ontic models do not provide an answer.  $\xi(m|M,\lambda)$  provides a probability and  $\lambda$  does little to explain the transition from potential to actual.
- Determinate ontic states do give an answer:

$$\xi(m|M,\lambda)\in\{0,1\}.$$

The MKC models are constructed precisely to show the logical possibility of a **non-contextual hidden variable theory**. Allegedly, this possibility was ruled out by the Kochen-Specker theorem.

## Assigning values to observables

Every non-degenerate self-adjoint operator A can be written as

$$A = a_1 |e_1\rangle \langle e_1| + \ldots + a_n |e_n\rangle \langle e_n|$$

with

$$\mathcal{B}_{A} = \{|e_{1}\rangle, \ldots, |e_{n}\rangle\}$$

an orthonormal basis.

- Assigning value a<sub>i</sub> to A
- = select vector  $|e_j\rangle$  from  $\mathcal{B}_A$ .

For every non-degenerate self-adjoint operator A and every function f, the observables f(A) and A can be measured jointly and the outcome for f(A) is  $f(a_i)$ .

• The definite value assigned to f(A) is determined by the definite value assigned to A.

## The Kochen-Specker Theorem

### Assigning definite values

An ontic state assigns definite values to observables by selecting for every orthonormal basis  $\mathcal{B}$  the "true" vector.

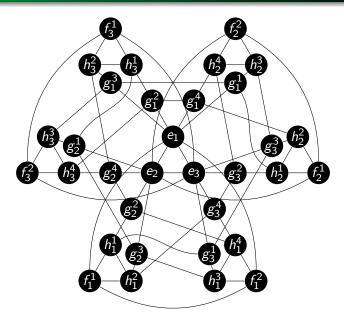
#### Non-Contextuality

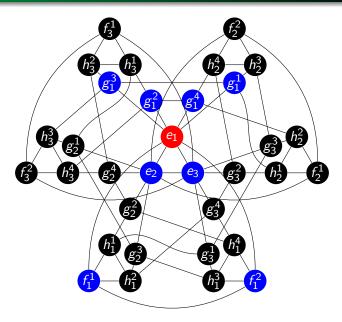
For every ontic state:  $|e\rangle$  is the "true" vector in an orthonormal basis  $\mathcal{B}$  iff it is the true vector in every other orthonormal basis  $\mathcal{B}'$  that contains  $|e\rangle$ .

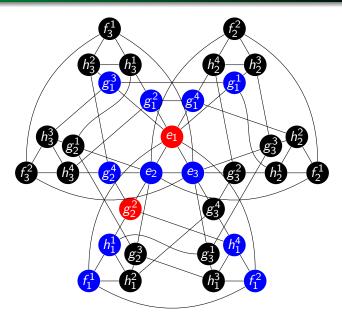
#### Kochen-Specker Theorem

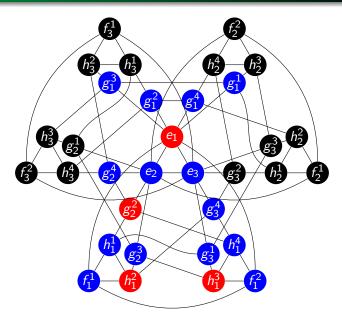
There is a finite set of orthonormal bases for which one cannot select true vectors in a non-contextual way.

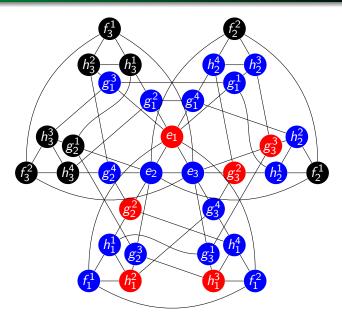
# Uncolorable graph in 3 dimensions (Peres-cube)









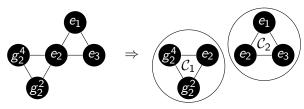


### Outline

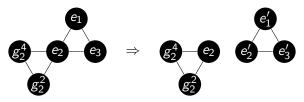
- The PBR Theorem
- ullet Two notions of  $\psi$ -onticity
- The KS Theorem
- The MKC models
- ullet  $\psi$ -onticity of the MKC models

## Circumventing the Kochen-Specker Theorem

Bell: Value assigned to an observable depends on the context  ${\mathcal C}$ 



MKC: Not every orthonormal basis represents an observable



but every orthonormal basis can be approximated by an observable

$$0<\|e_2-e_2'\|<\epsilon.$$

# The ontic models of Meyer, Kent and Clifton (MKC)

#### Definition

Two orthonormal bases

$$\mathcal{B}_{1}=\left\{ \left|e_{1}^{1}\right\rangle ,\ldots,\left|e_{n}^{1}\right\rangle \right\} ,\ \mathcal{B}_{2}=\left\{ \left|e_{1}^{2}\right\rangle ,\ldots,\left|e_{n}^{2}\right\rangle \right\}$$

are totally incompatible if  $\forall i, j = 1, \dots, n$  :

$$0<|\left\langle e_{i}^{1}\left|e_{j}^{2}\right\rangle \right|<1.$$

#### Trivial Theorem

Let  $\mathfrak{B}$  be any set of pairwise totally incompatible orthonormal bases. The set of all self-adjoint operators with eigenvectors in one of the bases in  $\mathfrak{B}$  is colorable.

### Non-Trivial Theorem (Clifton & Kent)

There exists a countable set  $\mathfrak B$  of pairwise totally incompatible orthonormal bases that lies dense in the set of all orthonormal bases.

# The ontic models of Meyer, Kent and Clifton (MKC)

- Ontic states:  $\Lambda = \{\lambda : \mathfrak{B} \to \{1, \dots, n\}\}.$
- ullet  $\Sigma=\sigma$ -algebra generated by cylinder sets.
- ullet  ${\cal P}$  determined by Born rule + independence of observables.
- $|e_i^k\rangle \in \mathcal{B}_k$ :

$$\begin{aligned} C_{e_i^k} &:= \left\{ \lambda \in \Lambda \mid \lambda(k) = i \right\}, \\ \mu_{\psi}(C_{e_i^k}) &:= |\langle \psi | e_i^k \rangle|^2. \end{aligned}$$

 $\bullet \ |e_{i_1}^{k_1}\rangle \in \mathcal{B}_{k_1}, \ |e_{i_2}^{k_2}\rangle \in \mathcal{B}_{k_2}, \ \ldots, \ |e_{i_n}^{k_n}\rangle \in \mathcal{B}_{k_n}:$ 

$$egin{aligned} C_{e_{i_1}^{k_1}, \dots, e_{i_n}^{k_n}} &:= igcap_{j=1}^n \left\{ \lambda \in \Lambda \, | \, \lambda(k_j) = i_j 
ight\}, \ \mu_{\psi}(C_{e_{i_1}^{k_1}, \dots, e_{i_n}^{k_n}}) &:= \prod_{j=1}^n |\langle \psi | e_j^k 
angle|^2. \end{aligned}$$

### Outline

- The PBR Theorem
- ullet Two notions of  $\psi$ -onticity
- The KS Theorem
- The MKC models
- ullet  $\psi$ -onticity of the MKC models

## Are the MKC models $\psi$ -ontic?



#### **PBR Theorem**

Any ontic model that reproduces the predictions of QM and satisfies the Preparation Independence Postulate is  $\psi$ -ontic.

The MKC models have never been properly defined for composite systems. But reasonable attempts violate PIP.

#### BCLM Theorem

Any ontic model that reproduces the predictions of QM is not maximally  $\psi$ -epistemic (but "almost"  $\psi$ -ontic).

## Are the MKC models $\psi$ -ontic?

#### $\psi$ -onticness:

For all pairs of quantum states  $\psi, \phi$ , for all corresponding probability measures  $\mu \in \Delta_{\psi}$ ,  $\nu \in \Delta_{\phi}$  in the ontic model, the variational distance  $D(\mu, \nu) := \sup_{\Omega \in \Sigma} |\mu(\Omega) - \nu(\Omega)|$  equals 1.

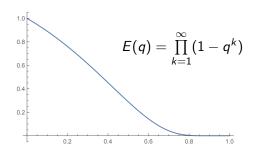
What is  $D(\mu_{\psi}, \mu_{\phi})$  in the MKC models?

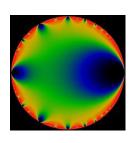
For cylinder sets:

$$\begin{split} &\mu_{\psi}\big(\mathit{C}_{e_{i_{1}}^{k_{1}},\ldots,e_{i_{n}}^{k_{n}}}\big) := \prod_{j=1}^{n} |\langle \psi | e_{i_{j}}^{k_{j}} \rangle|^{2} \rightarrow 1, \text{ as } n \rightarrow \infty, \\ &\mu_{\phi}\big(\mathit{C}_{e_{i_{1}}^{k_{1}},\ldots,e_{i_{n}}^{k_{n}}}\big) := \prod_{j=1}^{n} |\langle \phi | e_{i_{j}}^{k_{j}} \rangle|^{2} \rightarrow 0, \text{ as } n \rightarrow \infty. \end{split}$$

# Finding ontic states for $|\psi\rangle$

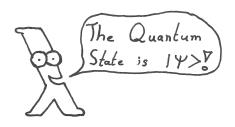
Wanted: 
$$\mu_{\psi}(\mathcal{C}_{e_{i_1}^{k_1},\dots,e_{i_n}^{k_n}}) := \prod_{j=1}^n |\langle \psi|e_{i_j}^{k_j} \rangle|^2 \to 1$$
, as  $n \to \infty$ ,





- Given  $\epsilon > 0$ , choose  $q_{\epsilon}$  such that  $E(x) > 1 \epsilon$ .
- ullet For every k choose  $(\mathcal{B}_{n_k},|e_{i_k}^{n_k}\rangle)$  such that  $|\langle\psi|e_{i_k}^{n_k}
  angle|^2>1-q_\epsilon^k$ .
- $\begin{array}{l} \bullet \ \, \mathsf{Set} \, \, \Lambda_\psi^\epsilon := \bigcap_{k=1}^\infty C_{\mathsf{e}_{i_k}^{n_k}}, \\ \\ \mathsf{then} \, \, \mu_\psi(\Lambda_\psi^\epsilon) > 1 \epsilon \, \, \mathsf{and} \, \, \mu_\phi(\Lambda_\psi^\epsilon) = 0 \, \, \mathsf{for all} \, \, |\phi\rangle. \\ \Rightarrow \, D(\mu_\psi, \mu_\phi) = 1 \\ \end{array}$

### How $\psi$ -ontic are the $\psi$ -ontic MKC models?



- There exists a function  $f: \Lambda \to \mathcal{P}$ , compatible with QM:  $\mu_{|\psi\rangle}(f^{-1}(|\psi\rangle)) = 1$ .
- ② If  $|\psi\rangle \neq |\phi\rangle$ , then  $\mu_{\psi}$  and  $\mu_{\phi}$  are non-overlapping.
  - The MKC-models are  $\psi$ -ontic in the second sense:  $D(\mu_{\psi}, \mu_{\phi}) = 1$ .
  - But we want the ontic state to tell us what  $|\psi\rangle$  is, i.e.,  $\Lambda_{\psi}\subset\Lambda$  such that

$$\mu_{\psi}(\Lambda_{\psi}) = 1, \ \mu_{\phi}(\Lambda_{\psi}) = 0$$

### How $\psi$ -ontic are the $\psi$ -ontic MKC models?



- $\Lambda_{\psi}^{\epsilon}$  does not contain all the  $\psi$ -ontic states:  $\mu_{\psi}(\Lambda_{\psi}^{\epsilon}) < 1$ .
- Moreover, there is no set  $\Lambda_{\psi}$  such that

$$\mu_{\psi}(\Lambda_{\psi}) = 1, \ \mu_{\phi}(\Lambda_{\psi}) = 0.$$

• Not all states in  $\Lambda_{\psi}^{\epsilon}$  are  $\psi$ -ontic states:

$$\Lambda_{\psi}^{\epsilon} \cap \Lambda_{\phi}^{\epsilon} \neq \varnothing$$
.

- The set  $\Lambda_{\psi}^{\epsilon}$  is underdetermined by  $(\psi, \epsilon)$ .
- More generally:  $f: \Lambda \to \mathcal{P}$  need not be unique, so we cannot speak of an identity relation.