

# $\psi$ -ontic models without $\psi$

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Where were you  
When the PBR-paper came out?

*"The wavefunction is a real physical object after all, say researchers"*

# Simpleminded counterexample?

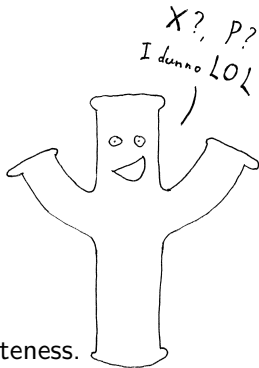
- Traditional hidden variables complement  $\psi$ :

$$(|\psi\rangle, \lambda)$$

- Why hidden variables? Incompleteness!

$$\Delta X \Delta P \geq \frac{\hbar}{2}$$

- But  $|\psi\rangle$  is a poor starting point for completeness.
- *Replace*  $|\psi\rangle$  with  $\lambda$ .  
Contextuality, non-locality, etc.



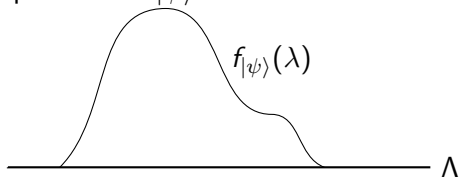
- 1 Recap of the PBR Theorem
- 2 Highlighting the important/problematic step
- 3 The counterexample
- 4 Conclusion

# Statement of the Theorem

## PBR Theorem

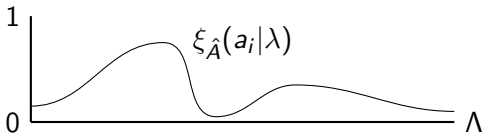
*Any ontic model that reproduces the predictions of QM and satisfies the Preparation Independence Postulate is  $\psi$ -ontic.*

- Preparation of  $|\psi\rangle$ :



$$\int f_{|\psi\rangle}(\lambda) d\lambda = 1$$

- Measurement of  $\hat{A}$ :



$$\sum_i \xi_{\hat{A}}(a_i|\lambda) = 1$$

- Compatibility:

$$\int \xi_{\hat{A}}(a_i|\lambda) f_{|\psi\rangle}(\lambda) d\lambda = |\langle \psi | a_i \rangle|^2$$

# Statement of the Theorem

## PBR Theorem

*Any ontic model that reproduces the predictions of QM and satisfies the Preparation Independence Postulate is  $\psi$ -ontic.*

*We call a hidden variable model  $\psi$ -ontic if every complete physical state or ontic state in the theory is consistent with only one pure quantum state; we call it  $\psi$ -epistemic if there exist ontic states that are consistent with more than one pure quantum state. – Harrigan, Spekkens 2010*

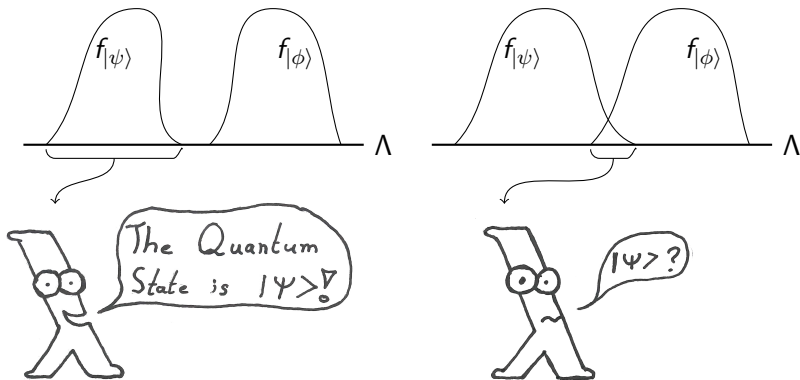
⇒ The ontic state  $\lambda$  determines the quantum state  $|\psi\rangle$ .

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The ontic state  $\lambda$  determines the quantum state  $|\psi\rangle$ :

*An important step towards the derivation of our result is the idea that the quantum state is physical if distinct quantum states correspond to non-overlapping distributions for  $\lambda$ . – Pusey, Barrett, Rudolph*



# Four levels of $\psi$ -onticity

- 1 There is a rule  $R$  assigning to every ontic state  $\lambda$  its corresponding quantum state  $R(\lambda)$  such that

$$\mathbb{P}_{|\psi\rangle}(\{\lambda \mid R(\lambda) = |\psi\rangle\}) = 1.$$

- 2 For every  $|\psi\rangle$  there is a set  $\Lambda_{|\psi\rangle}$  such that for all  $|\phi\rangle$ :

$$\mathbb{P}_{|\psi\rangle}(\Lambda_{|\psi\rangle}) = 1 \text{ and } \mathbb{P}_{|\phi\rangle}(\Lambda_{|\psi\rangle}) = 0.$$

- 3 For every  $|\psi\rangle$ ,  $\epsilon > 0$  there is a set  $\Lambda_{|\psi\rangle}^\epsilon$  such that for all  $|\phi\rangle$ :

$$\mathbb{P}_{|\psi\rangle}(\Lambda_{|\psi\rangle}^\epsilon) > 1 - \epsilon \text{ and } \mathbb{P}_{|\phi\rangle}(\Lambda_{|\psi\rangle}^\epsilon) < \epsilon.$$

- 4 For every  $|\psi\rangle$  and  $|\phi\rangle$  the distributions  $\mathbb{P}_{|\psi\rangle}, \mathbb{P}_{|\phi\rangle}$  are non-overlapping.

Toy examples for QM show

$$3 \not\Rightarrow 2 \not\Rightarrow 1$$

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# A contextual value definite ontic model

- A **context**  $C$  is a set of pairwise commuting 1-dimensional projectors that sum to  $\hat{1}$ .
- An **ontic state**  $\lambda$  is a function that assigns to each context a 1-dimensional projection operator such that

$$\lambda(C) \in C \quad \forall C.$$

- The **value of an observable**  $\hat{A}$  in the context  $C$  when the state is  $\lambda$  is

$$v_\lambda(\hat{A}|C) := \text{Tr}(\lambda(C)\hat{A}).$$

- Satisfies
  - $v_\lambda(\hat{A}|C) \in \sigma(\hat{A})$ ,
  - $v_\lambda(f(\hat{A})|C) = f(v_\lambda(\hat{A}|C))$ .

# Recovering the Born rule

- Set of ontic states:

$$\Lambda = \{\lambda : \mathcal{C} \rightarrow L(\mathcal{H}) \mid \lambda(C) \in C\}.$$

- $C_1, C_2, \dots, C_n$  contexts,
- $P_1, P_2, \dots, P_n$  projectors such that  $P_i \in C_i$ , then

$$\Delta_{C_1, \dots, C_n}^{P_1, \dots, P_n} = \{\lambda \in \Lambda \mid \lambda(C_i) = P_i\}.$$

- For any  $|\psi\rangle$ :

$$\mathbb{P}_{|\psi\rangle} \left( \Delta_{C_1, \dots, C_n}^{P_1, \dots, P_n} \right) = \prod_{i=1}^n \langle \psi | P_i | \psi \rangle.$$

The model is  $\psi$ -ontic

For every  $|\psi\rangle$  there is a set  $\Lambda_{|\psi\rangle}$  such that for all  $|\phi\rangle$ :

$$\mathbb{P}_{|\psi\rangle}(\Lambda_{|\psi\rangle}) = 1 \text{ and } \mathbb{P}_{|\phi\rangle}(\Lambda_{|\psi\rangle}) = 0.$$

*Proof:*

- Choose countable sequence  $C_1, C_2, \dots$  such that

$$|\psi\rangle\langle\psi| \in C_i \quad \forall i$$

- Define

$$\Lambda_{|\psi\rangle} = \{\lambda \in \Lambda \mid \lambda(C_i) = |\psi\rangle\langle\psi|\}$$

But without  $\psi$

If  $\langle\phi|\psi\rangle \neq 0$ , then

$$\Lambda_{|\psi\rangle} \cap \Lambda_{|\phi\rangle} \neq \emptyset.$$

Negative conclusion:

- In the ontic model there are ontic states that are consistent with more than one pure quantum state. So, according to Harrigan & Spekkens, it is  $\psi$ -epistemic.

⇒  $\psi$ -ontology theorems fails!

Positive conclusion:

*I propose that we stop talking about the ill-defined notion of quantum state realism, and that we start talking instead about these sorts of question – e.g. whether quantum theory comes with objective standards for the ascription of states to physical situations. – Halvorson 2018*

- The fact that quantum states correspond to non-overlapping distributions indicates that two agents using distinct quantum states are substantially disagreeing about *something*.